

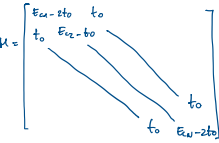
In  $i$ th point, the Schrödinger is discretized as:

$$\frac{1}{2} \psi_i = -\frac{\hbar^2}{2m\Delta x^2} (\psi_{i+1} - 2\psi_i + \psi_{i-1}) + E \psi_i$$

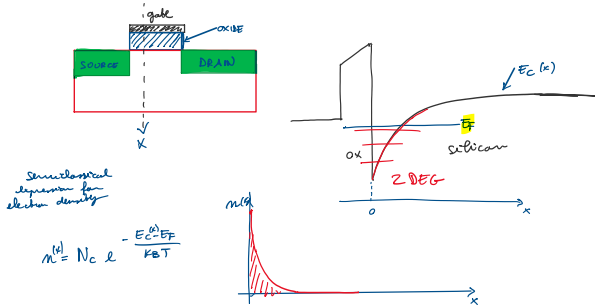
$$\psi_{i+1} + (E_{cc} - 2t_0) \psi_i + t_0 \psi_{i-1} = E \psi_i$$

$$[H][\psi] = E[\psi]$$

$$\psi_0 = 0 \wedge \psi_{N+1} = 0$$



this is like considering  $E_{c0} \rightarrow +\infty$  and  $E_{cN+1} \rightarrow +\infty$



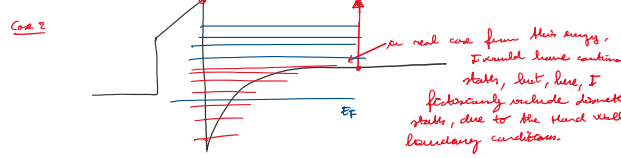
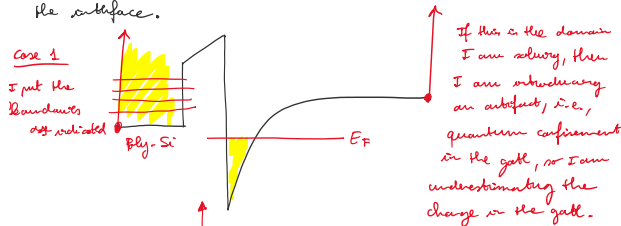
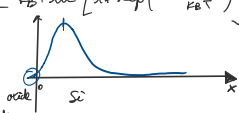
Semiclassical expression for electron density

$$n^{(H)} = N_c \exp\left(-\frac{E_c - E_F}{k_B T}\right)$$

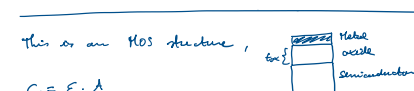
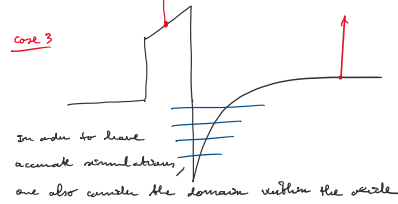
In the case of quantum confinement:

$$N_{2DEG} = \sum_i |\psi_i|^2 \frac{\sqrt{m_0^* m_0}}{\pi \hbar^2} k_B T \ln \left[ 1 + \exp\left(-\frac{E_i - E_F}{k_B T}\right) \right]$$

$\psi_i$  is an Airy function  
 $\psi_i$  matches the size, and the maximum is not at the interface.



This case completely neglects the penetration of the  $\psi$  in the oxide  $\Rightarrow$  this could introduce a large error



This is an MOS structure,  $C = \epsilon \frac{A}{t_{ox}}$ , because the charge is squeezed at the interface.

In quantum confinement case, the centroid of charge is not at  $x=0$  but it is, right shifted

$$t_a = \frac{\int x n(x) dx}{\int n(x) dx} \quad C \approx \epsilon \frac{A}{t_{ox} + t_a} \quad \frac{1}{C} = \frac{t_{ox}}{\epsilon} + \frac{t_a}{\epsilon}$$

QM is also leading to an increase of the threshold voltage of the MOS, due to fewer available states as compared to the semiclassical case

